"Ideas, like ghosts, must be spoken to a little before they will explain themselves"

Charles Dickens, Dombey and Son

Dynamics of Fluid Loaded Structures (with application to Loose Parts Monitoring (LPM) in nuclear steam plants)

The second li

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<u>Abstract</u> Parts that come loose in the cooling circuit can endanger the integrity of a nuclear steam plant. Systems to detect the impacts that occur are installed in the plants but they face difficulties in the detection and classification of such impacts because of the noisy vibrations induced by turbulent flow. Optimizing the performance of these LPM systems depends on knowing the "signature" of the induced vibration signals and these signals are affected by the structural acoustics of the plant components that are internally loaded by water. This presentation reviews the acoustics of fluid loaded plate structures and the dynamics of impact collisions for such structures.

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Boiling Water Reactor



Pressurized Water Reactor

Impact dynamics





Assumptions made in the structural analysis

- The structure can be modeled as an isotropic flat plate. This primarily means that the stiffening effect of curvature is ignored
- The structural waves of importance are bending waves Although in-plane dilatation and shear waves are generated, they likely store little energy and have small effect on the measurements of vibration.
- The fluid is assumed incompressible. This should be acceptable as long as the phase speed of the bending waves is small compared to the speed of sound in the fluid.
- The fluid is assumed inviscid. That means that the small amount of damping provided by viscous losses at the wall/fluid interface is incorporated as part of the structural damping.
- Bouncing of the part after collision with the wall does not occur. This is done for convenience, it would be possible to include bouncing in a more elaborate analysis.

Equations governing the structural dynamics of fluid loaded flat plate (phase and group speed and modal density)

$$(EI\nabla^4 - \omega^2 \rho_s)y = EI(k^4 - k_b^4)y = -\omega^2 \rho_f y/k$$

k=*ω*/*c*, *kb*=*ω*/*cb*

$$\frac{\rho_f}{\omega\rho_s}c^5 + c^4 - c_b^4 = 0$$

$$c_g = c(1 - d \, \ell n \, c \, / \, d \, \ell n \, \omega)^{-1}$$

$$n(\omega) = \omega A_p / 2\pi c c_g$$





Conductance and Susceptance

$$G(\omega) = \pi n(\omega) / 2M(\omega)$$



 $n(\omega) \sim \omega p$ $Y(\omega) = G(\omega)(1 - j2R / \pi);$

$$R = 2p/(p^2 - 1)$$





Transfer functions

$$\langle f^2 \rangle G(\omega) = \omega \eta M(\omega) \langle v^2 \rangle$$









e(t)=M/2
e(t)=Eexp(-ωηt)
S(ω)=½M
$$\int <|v|^2 > dt = (\frac{1}{2}M/\omega^2) \int <|a|^2 > dt = E/ωη$$

E(ω)=ωηS(ω)









Samples of impact forces and resulting acceleration

Modal parameters as noise to be rejected



 $x(t)=sin(2\pi f_o t)exp(-\beta_o t)$

parameters: f_o , β_o

 $h(t) = \sum_{m} \psi_{m}^{s} \psi_{m}^{o} \sin(2\pi f_{m} t) \exp(-\beta_{m} t) \quad m = 1, 2, \dots$ parameters: $f_{m}, \beta_{m}, \psi_{m}^{s}, \psi_{m}^{o} m = 1, 2, \dots$

Therefore

y(t): parameters: f_o , β_o ; f_m , β_m , ψ_m^s , ψ_m^o m=1,2,...

Question:

How to get rid of f_m , β_m , ψ_m^{s} , ψ_m^{o} m=1,2,... and retrieve f_o , β_o ?



Matched Filtering



Figure 1. Matched filter applied to simple pulse waveform









