

Problem 1

$$i) \quad X(f) = \begin{cases} jX_0 & f \in [f_0, f_0+w] \\ -jX_0 & f \in [-f_0-w, -f_0] \end{cases}$$

Note that conjugate symmetry holds ($X(f) = X^*(-f)$)

Obviously, $\tilde{X}(f) = jX_0$, $f \in [0, w]$ (carrier freq = f_0)

$$X_c(f) \triangleq \text{Even}[\tilde{X}(f)] = \frac{1}{2} (\tilde{X}(f) + \tilde{X}^*(-f)) = \begin{cases} j\frac{X_0}{2}, & f \in [0, w] \\ -j\frac{X_0}{2}, & f \in [-w, 0] \end{cases}$$

$$X_s(f) \triangleq j \text{ odd } \tilde{X}(f) = \frac{j}{2} (\tilde{X}(f) - \tilde{X}^*(-f)) = \begin{cases} -X_0/2, & f \in [0, w] \\ X_0/2, & f \in [-w, 0] \end{cases}$$

$$\begin{aligned} ii) \quad X_c(t) &\triangleq \text{Re}[\tilde{X}(t)] = \text{Re}\left[\int_0^w jX_0 e^{j2\pi ft} df\right] = \\ &= \text{Re}\left[jX_0 \frac{e^{j2\pi wt} - 1}{j2\pi t}\right] = \text{Re}\left[X_0 e^{j2\pi \frac{w}{2}t} \frac{\sin(\pi wt)}{\pi t}\right] \\ &= X_0 \text{sinc}(\pi wt) \text{Re}\left[e^{j2\pi \frac{w}{2}t}\right] = X_0 \text{sinc}(\pi wt) \cos\left(2\pi \frac{w}{2}t\right) \end{aligned}$$

$$\begin{aligned} X_s(t) &\triangleq -\text{Im}[\tilde{X}(t)] = -\text{Im}\left[\int_0^w jX_0 e^{j2\pi ft} df\right] \\ &= -\text{Im}\left[X_0 e^{j2\pi \frac{w}{2}t} \text{sinc}(\pi wt)\right] = -X_0 \text{sinc}(\pi wt) \sin\left(2\pi \frac{w}{2}t\right) \end{aligned}$$

$$iii) \quad X_1(f) = \begin{cases} jX_0 & , f \in [f_0 + \delta f, f_0 + W + \delta f] \\ -jX_0 & , f \in [-f_0 - W - \delta f, -f_0 - \delta f] \end{cases}$$

thus $\tilde{X}_1(f) = jX_0 \quad f \in [\delta f, \delta f + W]$ or
 $\tilde{X}_1(f) = \tilde{X}(f - \delta f)$

$$X_{1e}(f) = \text{Even}[\tilde{X}_1(f)] = \begin{cases} jX_0/2 & , f \in [\delta f, W + \delta f] \\ -jX_0/2 & , f \in [-W - \delta f, -\delta f] \end{cases}$$

$$X_{1s}(f) = \text{odd}[\tilde{X}_1(f)] = \begin{cases} -X_0/2 & , f \in [\delta f, W + \delta f] \\ -X_0/2 & , f \in [-W - \delta f, -\delta f] \end{cases}$$

$$x_{1e}(t) = \text{Re}[\tilde{x}_1(t)] = \text{Re}\left[\int_{\delta f}^{\delta f + W} jX_0 e^{j2\pi ft} df\right] =$$

$$= \text{Re}\left[e^{+j2\pi\delta ft} \int_0^W jX_0 e^{j2\pi ft} df\right] = \text{Re}\left[e^{+j2\pi\delta ft} X_0 e^{j2\pi\frac{W}{2}t} \text{sinc}(nWt)\right]$$

$$= X_0 \text{sinc}(nWt) \cos\left(2\pi\left(\frac{W}{2} + \delta f\right)t\right)$$

similarly, $x_{1s}(t) = -X_0 \text{sinc}(nWt) \sin\left(2\pi\left(\frac{W}{2} + \delta f\right)t\right)$

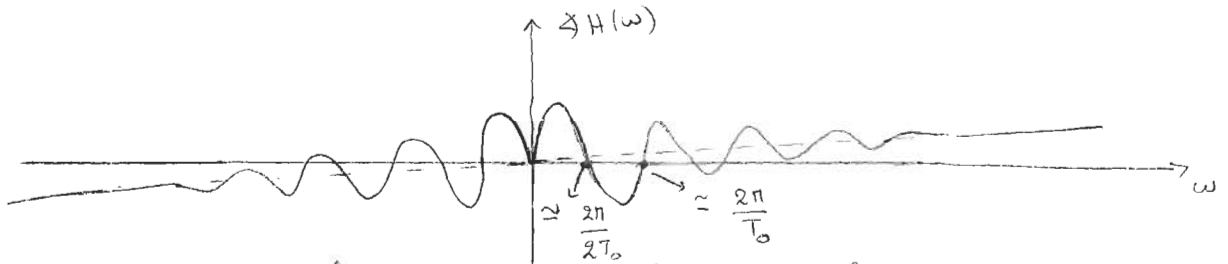
iv) if X_0 is complex then from part ii)

$$x_c(t) = \text{Re}\left[X_0 e^{j2\pi\frac{W}{2}t} \text{sinc}(nWt)\right] = \text{sinc}(nWt) \text{Re}\left[X_0 e^{j2\pi\frac{W}{2}t}\right]$$

$$x_s(t) = -\text{sinc}(nWt) \text{Im}\left[X_0 e^{j2\pi\frac{W}{2}t}\right]$$

Problem 2

i) $\hat{H}(\omega) = (s_0 \omega + \alpha \sin(|\omega T_0|) e^{-|\omega T_1|}) R$



the contribution of the first nulls comes from $\sin(|\omega T_0|)$ only because $T_0 \gg T_1$ (exp. decay does not affect them) and $\alpha \gg s_0 \omega$

ii) $T_g(\omega) \triangleq -\frac{1}{\omega} \hat{H}(\omega) = -R \left(s_0 + \alpha \sin(|\omega T_0|) \frac{e^{-|\omega T_1|}}{\omega} \right)$

I need to know $\frac{d}{dx} |x| = \frac{d}{dx} \sqrt{x^2} = \frac{1}{2\sqrt{x^2}} 2x = \frac{x}{|x|}$

$= \text{sgn}(x) \frac{|x|}{|x|} = \text{sgn}(x) = \begin{cases} 1, & x > 0 \\ -1, & x < 0 \end{cases}$

hence, $T_g(\omega) \triangleq -\frac{d}{d\omega} \hat{H}(\omega) =$

$= -R \left[s_0 + \alpha \left(\cos(|\omega T_0|) \text{sgn}(\omega T_0) T_0 e^{-|\omega T_1|} + \sin(|\omega T_0|) e^{-|\omega T_1|} \frac{1}{T_1} (-\text{sgn}(\omega T_1)) \right) \right]$

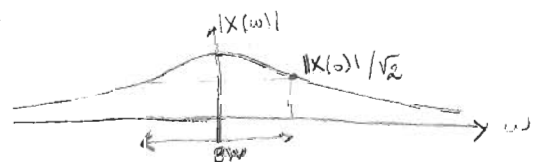
$= -R \left[s_0 + \alpha e^{-|\omega T_1|} \left(T_0 \text{sgn}(\omega T_0) \cos(|\omega T_0|) - T_1 \text{sgn}(\omega T_1) \sin(|\omega T_0|) \right) \right]$

for $\omega > 0$ $T_g(\omega) = -R \left[s_0 + \alpha e^{-\omega T_1} \left(T_0 \cos(\omega T_0) - T_1 \sin(\omega T_0) \right) \right]$

iii) $x(t) = x_0 e^{-2t/T_0} u(t)$ input
 $X(f) = \int_0^{\infty} x_0 e^{-2t/T_0} e^{-j2\pi ft} dt = x_0 \int_0^{\infty} e^{-j2\pi t(f + \frac{2}{T_0 2\pi j})} dt$

$x_0 \left(\frac{0 - 1}{-j2\pi(f + \frac{2}{T_0 2\pi j})} \right) = \frac{x_0}{j2\pi f + \frac{2}{T_0}}$ $\Rightarrow X(\omega) = \frac{x_0}{j\omega + \frac{2}{T_0}}$

$\Rightarrow |X(\omega)| = \frac{|x_0|}{\sqrt{\omega^2 + 4/T_0^2}}$



Define BW where the magnitude of $X(\omega)$ falls to $\frac{|X(0)|}{\sqrt{2}}$. Hence $\frac{|X(0)|}{\sqrt{2}} = \frac{|X_0|}{\sqrt{\omega^2 + 4/T_0^2}} \Rightarrow$

$$\frac{|X_0|}{\sqrt{2} \frac{2}{T_0}} = \frac{|X_0|}{\sqrt{\omega^2 + 4/T_0^2}} \Rightarrow \omega^2 + \frac{4}{T_0^2} = 2 \frac{4}{T_0^2} \Rightarrow$$

$$\omega = \pm \frac{2}{T_0} \quad \text{thus BW of } x(t) \text{ is } \pm \frac{2}{T_0}.$$

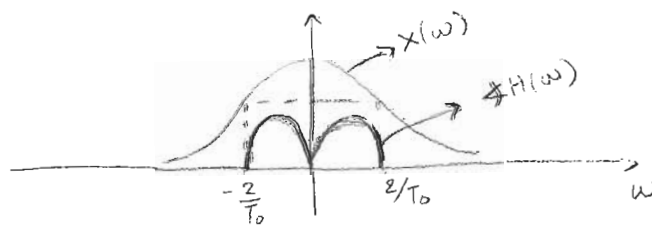
The output of the filter is:

$$Y(\omega) = X(\omega) H(\omega) = X(\omega) |H(\omega)| e^{j \angle H(\omega)} =$$

$$= \frac{X_0}{j\omega + \frac{2}{T_0}} |H(\omega)| e^{jR(\alpha_0 \omega + \alpha \sin |\omega T_0|) e^{-|\omega T_0|}}$$

The region of interest is when $\omega \in \left[-\frac{2}{T_0}, \frac{2}{T_0}\right]$ i.e. the "effective" BW of $x(t)$.

In this region the contribution of $\angle H(\omega)$ comes from $\alpha \sin |\omega T_0|$ only so we can approximate $\angle H(\omega) \approx \alpha R \sin |\omega T_0|$.



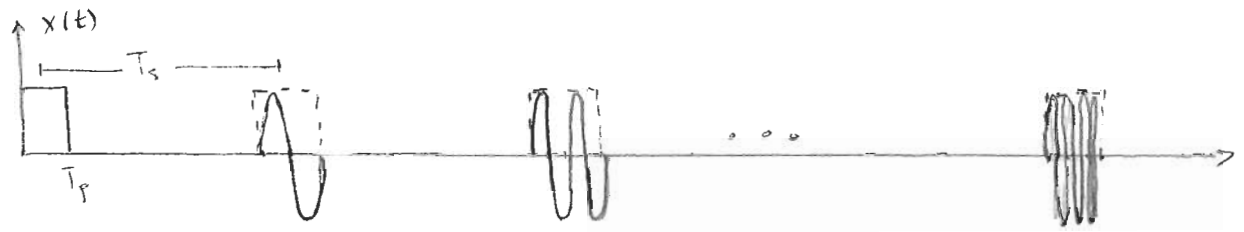
Obviously dispersion cannot be neglected so

$$Y(\omega) \approx \frac{X_0}{j\omega + \frac{2}{T_0}} H_0 e^{jR \alpha \sin |\omega T_0|}$$

$$\text{and } y(t) \approx \int_{-2/T_0}^{2/T_0} \frac{X_0 H_0}{j\omega + \frac{2}{T_0}} e^{jR \alpha \sin |\omega T_0|} e^{j\omega t} \frac{d\omega}{2\pi}$$

Problem 3

i) The transmitted signal is:



We want $\int x^2(t) dt = 1 \Rightarrow$

$$\sum_{n=1}^N \int_{(n-1)T_s}^{(n-1)T_s + T_p} A^2 \sin^2(\omega_n(t - (n-1)T_s)) dt = 1 \Rightarrow$$

$$\sum_{n=1}^N \int_0^{T_p} A^2 \sin^2(\omega_n t) dt = 1 \Rightarrow \sum_{n=1}^N \int_0^{T_p} A^2 \frac{(1 - \cos(2\omega_n t))}{2} dt = 1$$

$$\frac{A^2}{2} \sum_{n=1}^N \int_0^{T_p} dt - \int_0^{T_p} \cos(2 \cdot 2\pi(n-1) \frac{1}{T_p} t) dt = 1 \Rightarrow$$

$$\frac{A^2 N T_p}{2} = 1 \Rightarrow A = \left(\frac{2}{N T_p} \right)^{1/2}$$

ii) $\phi(0, \Delta f) = \int_{-\infty}^{\infty} x^2(t) e^{j2\pi \Delta f t} dt =$

$$= \sum_{n=1}^N \int_{(n-1)T_s}^{(n-1)T_s + T_p} A^2 \sin^2(\omega_n(t - (n-1)T_s)) e^{j2\pi \Delta f t} dt =$$

$$= \sum_{n=1}^N \int_0^{T_p} A^2 \sin^2(\omega_n t) e^{j2\pi \Delta f (t + (n-1)T_s)} dt =$$

$$= \sum_{n=1}^N e^{-j2\pi \Delta f (n-1)T_s} \int_0^{T_p} \frac{A^2}{2} (1 - \cos(2\omega_n t)) e^{j2\pi \Delta f t} dt$$

$$= \sum_{n=1}^N \frac{A^2}{2} e^{-j2\pi \Delta f (n-1)T_s} \left[\int_0^{T_p} 1 \cdot e^{-j2\pi \Delta f t} dt - \int_0^{T_p} \cos(2\omega_n t) e^{-j2\pi \Delta f t} dt \right]$$

$$\int_0^{T_p} e^{-j2\pi \Delta f t} dt = F \left[\text{rect}_{T_p} \right] = T_p e^{-j2\pi \Delta f T_p/2} \text{sinc}(\pi T_p \Delta f)$$

$$\int_0^{T_p} \cos(2\omega_n t) e^{-j2\pi \Delta f t} dt = F \left[\text{rect}_{T_p} \cdot \cos(2\omega_n t) \right]$$

$$= T_p e^{-j2\pi \Delta f T_p/2} \text{sinc}(\pi T_p \Delta f) * \frac{\uparrow^{1/2} \quad \uparrow^{1/2}}{\frac{-2(n-1)}{T_p} \quad \frac{2(n-1)}{T_p}} \Delta f$$

$$= \frac{T_p}{2} e^{-j2\pi (\Delta f \pm \frac{(n-1)\omega}{T_p}) \frac{T_p}{2}} \text{sinc} \left(\pi T_p \left(\Delta f \pm \frac{2(n-1)\omega}{T_p} \right) \right)$$

Hence the end result is:

$$\phi(o, \Delta f) = A^2 \frac{1}{2} \sum_{n=1}^N e^{-j2\pi \Delta f (n-1)T_s} \left[T_p e^{-j2\pi \Delta f T_p/2} \text{sinc}(\pi T_p \Delta f) - \frac{T_p}{2} e^{-j2\pi (\Delta f + \frac{(n-1)\omega}{T_p}) \frac{T_p}{2}} \text{sinc} \left(\pi T_p \left(\Delta f + \frac{2(n-1)\omega}{T_p} \right) \right) - \frac{T_p}{2} e^{-j2\pi (\Delta f - \frac{(n-1)\omega}{T_p}) \frac{T_p}{2}} \text{sinc} \left(\pi T_p \left(\Delta f - \frac{2(n-1)\omega}{T_p} \right) \right) \right]$$

For Doppler accuracy the shifted $\text{sinc} \left(\pi T_p \left(\Delta f - \frac{2(n-1)\omega}{T_p} \right) \right)$ do not contribute so we approximate.

$$\begin{aligned} \phi(o, \Delta f) &\approx \frac{A^2}{2} \sum_{n=1}^N \left[e^{-j2\pi \Delta f (n-1)T_s} T_p e^{-j2\pi \Delta f T_p/2} \text{sinc}(\pi T_p \Delta f) \right] \\ &- T_p e^{-j2\pi \Delta f T_p/2} \text{sinc}(\pi T_p \Delta f) = \\ &= A^2 \frac{T_p}{2} \sum_{n=2}^N e^{-j2\pi \Delta f (n-1)T_s} \cdot e^{-j2\pi \Delta f T_p/2} \text{sinc}(\pi T_p \Delta f) \end{aligned}$$

$$= A^2 \frac{T_p}{2} \text{sinc}(n T_p \Delta f) e^{-j n \Delta f T_p} \underbrace{\sum_{n=2}^N e^{-j n \Delta f (n-1) T_s}}_{\text{sum of finite geometric series}}$$

$$\sum_{k=1}^N r^k = \frac{r(1-r^N)}{1-r}$$

$$= A^2 \frac{T_p}{2} \text{sinc}(n T_p \Delta f) e^{-j n \Delta f T_p} e^{-j 2 n \Delta f T_s} \frac{(1 - e^{-j 2 n \Delta f (N-1) T_s})}{1 - e^{-j 2 n \Delta f T_s}}$$

$$= A^2 \frac{T_p}{2} e^{-j n \Delta f (T_p + T_s)} \text{sinc}(n T_p \Delta f) \frac{\sin(n(N-1) T_s \Delta f)}{\sin(n T_s \Delta f)} \frac{e^{-j n \Delta f (n-1) T_s}}{e^{-j n \Delta f T_s}}$$

So $\phi(0, \Delta f)$ has peaks of width $\frac{1}{(N-1)T_s}$ and is periodic at $\Delta f = \frac{1}{T_s}$. This implies that the Doppler resolution is determined by the signal duration but there are ambiguous sidelobes at $\frac{1}{T_s}$.

Observing the structure of $\phi(0, \Delta f)$ there exist additional sidelobes @ $\frac{n}{T_p}$ $n=1, \dots, N$.

$$\phi(\Delta T, 0) = \int_{-\infty}^{\infty} x(t - \frac{\Delta T}{2}) x(t + \frac{\Delta T}{2}) dt = \int_{-\infty}^{\infty} x(t) x(t - \Delta T) dt$$

when $0 \leq \Delta T \leq T_p$ we are looking at the mainlobe width

$$\phi(\Delta T, 0) = \sum_{n=1}^N \int_{(n-1)T_s + \Delta T}^{(n-1)T_s + T_p} A \sin(\omega_n(t - (n-1)T_s)) A \sin(\omega_n(t - (n-1)T_s - \Delta T)) dt$$

using $\sin A \sin B = \frac{1}{2} (\cos(A-B) - \cos(A+B))$

$$\phi(\Delta T, 0) = \sum_{n=1}^N \frac{A^2}{2} \int_{(n-1)T_s + \Delta T}^{(n-1)T_s + T_p} \cos(\omega_n \Delta T) - \cos(2\omega_n(t - (n-1)T_s) - \omega_n \Delta T) dt$$

$$= \sum_{n=1}^N \frac{A^2}{2} \left[\cos(\omega_n \Delta T) (T_p - \Delta T) - \int_{\Delta T}^{T_p} \cos(2\omega_n t - \omega_n \Delta T) dt \right]$$

$$= \sum_{n=1}^N \frac{A^2}{2} \left[(T_p - \Delta T) \cos(\omega_n \Delta T) - \frac{\sin \omega_n (2T_p - \Delta T) - \sin(\omega_n \Delta T)}{2\omega_n} \right]$$

Observe that the first term of the sum defines the range resolution which is $\Delta T = T_p$ (first zero). Hence range resolution $\sim T_p$.

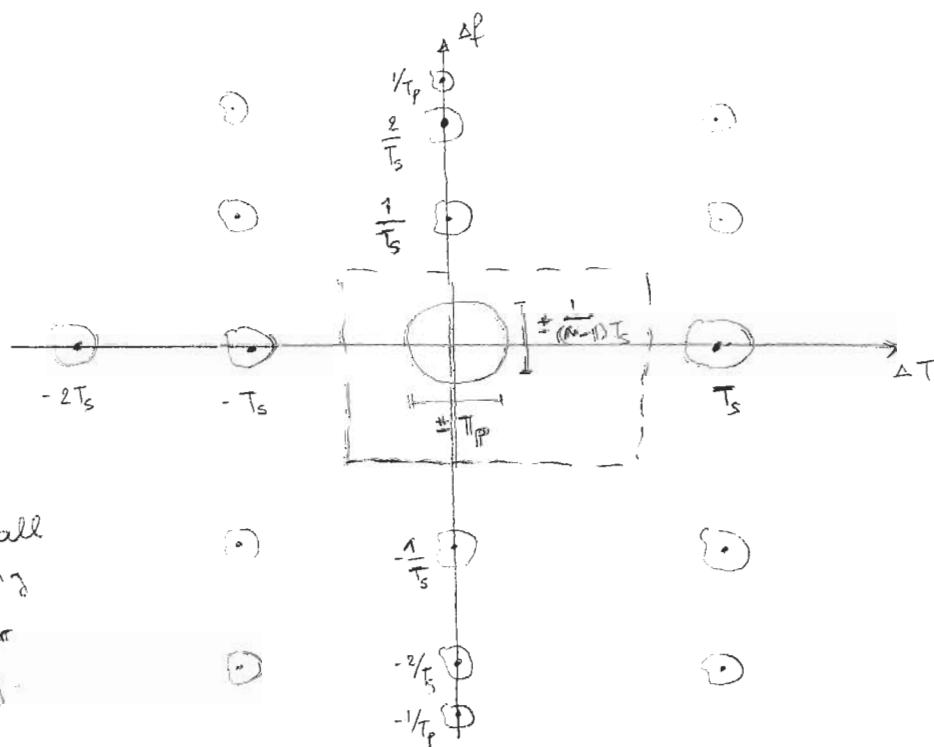
When $T_p < \Delta T < T_s - T_p$ then there is no overlap so $\phi(\Delta T, 0) = 0$.

When $T_s \leq \Delta T \leq T_s + T_p$ then we are looking at the width of the first sidelobe

$$\begin{aligned} \phi(\Delta T, 0) &= \sum_{n=2}^N \int_{(n-2)T_s + \Delta T}^{(n-1)T_s + T_p} A \sin \omega_n(t - (n-1)T_s) A \sin \omega_{n-1}(t - (n-2)T_s - \Delta T) dt \\ &= \sum_{n=2}^N \frac{A^2}{2} \int_{(n-2)T_s + \Delta T}^{(n-1)T_s + T_p} \left[\cos \left[\omega_n(t - (n-1)T_s) - \omega_{n-1}(t - (n-2)T_s - \Delta T) \right] - \right. \\ &\quad \left. \cos \left[\omega_n(t - (n-1)T_s) + \omega_{n-1}(t - (n-2)T_s - \Delta T) \right] \right] dt \end{aligned}$$

due to this constructive and destructive summation of cosines the sidelobe around T_s is severely smaller than the mainlobe.

So how does $\phi(\Delta T, \Delta f)$ look like?



make T_p small for range gating
 increase N for Doppler gating.